

Progress of the AMIDAS Package for Reconstructing WIMP Properties

Chung-Lin Shan

Xinjiang Astronomical Observatory
Chinese Academy of Sciences

4th International Workshop on
Dark Matter, Dark Energy, and Matter-Antimatter Asymmetry
December 31, 2016



AMIDAS package

Motivation

Reconstruction of the 1-D WIMP velocity distribution

With measured recoil energies

With a non-negligible threshold energy

With a model of the WIMP velocity distribution

Determination of the WIMP mass

Determinations of the WIMP-nucleon couplings

Estimation of the SI scalar WIMP-nucleon coupling

Determinations of ratios of WIMP-nucleon cross sections

Summary



AMIDAS package

– A Model-Independent Data Analysis System



AMIDAS package

- AMIDAS: A Model-Independent Data Analysis System
for direct Dark Matter detection experiments and phenomenology



AMIDAS package

- AMIDAS: A Model-Independent Data Analysis System
for direct Dark Matter detection experiments and phenomenology
 - DAMNED Dark Matter Web Tool (ILIAS Project)
<http://pisrv0.pit.physik.uni-tuebingen.de/darkmatter/amidas/>
[CLS, AIP Conf. Proc. 1200, 1031; arXiv:0910.1971; Phys. Dark Univ. 5-6, 240 (2014)]
 - TiResearch (Taiwan interactive Research)
<http://www.tir.tw/phys/hep/dm/amidas/>



AMIDAS package

- AMIDAS: A Model-Independent Data Analysis System
for direct Dark Matter detection experiments and phenomenology
 - DAMNED Dark Matter Web Tool (ILIAS Project)
<http://pisrv0.pit.physik.uni-tuebingen.de/darkmatter/amidas/>
[CLS, AIP Conf. Proc. 1200, 1031; arXiv:0910.1971; Phys. Dark Univ. 5-6, 240 (2014)]
 - TiResearch (Taiwan interactive Research)
<http://www.tir.tw/phys/hep/dm/amidas/>
 - Online interactive simulation/data analysis system




AMIDAS package

- AMIDAS: A Model-Independent Data Analysis System for direct Dark Matter detection experiments and phenomenology
 - DAMNED Dark Matter Web Tool (ILIAS Project)
<http://pisrv0.pit.physik.uni-tuebingen.de/darkmatter/amidas/>
[CLS, AIP Conf. Proc. 1200, 1031; arXiv:0910.1971; Phys. Dark Univ. 5-6, 240 (2014)]
 - TiResearch (Taiwan interactive Research)
<http://www.tir.tw/phys/hep/dm/amidas/>
 - Online interactive simulation/data analysis system
 - Full Monte Carlo simulations
 - Theoretical estimations
 - Real/pseudo- data analyses


AMIDAS package

- AMIDAS: A Model-Independent Data Analysis System for direct Dark Matter detection experiments and phenomenology

AMIDAS: A Model-Independent Data Analysis System for Direct Dark Matter Detection Experiments and Phenomenology



Chung-Lin Shan



[Dark Matter Online Tools](#)
[DAMNED](#)
AMIDAS: [restart](#), [new window](#), [last results](#)

Established: January 11, 2009
Last upgraded: March 18, 2015

List of AMIDAS functions

Choose one of the following functions

- Generation of WIMP signals with/without background events
- (Bayesian) reconstruction of the one-dimensional velocity distribution function of halo WIMPs ^{new}
- Determination of the WIMP mass
- Estimation of the spin-independent (SI) WIMP-nucleon coupling
- Determinations of ratios between different WIMP-nucleon couplings/cross sections



Motivation

Motivation

- Differential event rate for elastic WIMP-nucleus scattering

$$\frac{dR}{dQ} = \mathcal{A} F^2(Q) \int_{v_{\min}(Q)}^{v_{\max}} \left[\frac{f_1(v)}{v} \right] dv$$

Here

$$v_{\min}(Q) = \alpha \sqrt{Q}$$

is the minimal incoming velocity of incident WIMPs that can deposit the recoil energy Q in the detector,

$$\mathcal{A} \equiv \frac{\rho_0 \sigma_0}{2m_\chi m_{r,N}^2} \quad \alpha \equiv \sqrt{\frac{m_N}{2m_{r,N}^2}} \quad m_{r,N} = \frac{m_\chi m_N}{m_\chi + m_N}$$

ρ_0 : WIMP density near the Earth

σ_0 : total cross section ignoring the form factor suppression

$F(Q)$: elastic nuclear form factor

$f_1(v)$: one-dimensional velocity distribution of halo WIMPs

Motivation

- Differential event rate for elastic WIMP-nucleus scattering

$$\frac{dR}{dQ} = \mathcal{A} F^2(Q) \int_{v_{\min}(Q)}^{v_{\max}} \left[\frac{f_1(v)}{v} \right] dv$$

Here

$$v_{\min}(Q) = \alpha \sqrt{Q}$$

is the minimal incoming velocity of incident WIMPs that can deposit the recoil energy Q in the detector,

$$\mathcal{A} \equiv \frac{\rho_0 \sigma_0}{2m_\chi m_{r,N}^2} \quad \alpha \equiv \sqrt{\frac{m_N}{2m_{r,N}^2}} \quad m_{r,N} = \frac{m_\chi m_N}{m_\chi + m_N}$$

Particle physics

ρ_0 : WIMP density near the Earth

σ_0 : total cross section ignoring the form factor suppression

$F(Q)$: elastic nuclear form factor

$f_1(v)$: one-dimensional velocity distribution of halo WIMPs

Motivation

- Differential event rate for elastic WIMP-nucleus scattering

$$\left(\frac{dR}{dQ}\right) = \mathcal{A} F^2(Q) \int_{v_{\min}(Q)}^{v_{\max}} \left[\frac{f_1(v)}{v}\right] dv$$

Here

$$v_{\min}(Q) = \alpha \sqrt{Q}$$

is the minimal incoming velocity of incident WIMPs that can deposit the recoil energy Q in the detector,

$$\mathcal{A} \equiv \frac{\rho_0 \sigma_0}{2m_\chi m_{r,N}^2} \quad \alpha \equiv \sqrt{\frac{m_N}{2m_{r,N}^2}} \quad m_{r,N} = \frac{m_\chi m_N}{m_\chi + m_N}$$

ρ_0 : WIMP density near the Earth

σ_0 : total cross section ignoring the form factor suppression

$F(Q)$: elastic nuclear form factor

$f_1(v)$: one-dimensional velocity distribution of halo WIMPs

Astrophysics

Motivation

- Differential event rate for elastic WIMP-nucleus scattering

$$\underbrace{\left(\frac{dR}{dQ}\right)}_{\text{green arrow}} = \mathcal{A} F^2(Q) \int_{v_{\min}(Q)}^{v_{\max}} \left[\frac{f_1(v)}{v} \right] dv$$

Here

$$v_{\min}(Q) = \alpha \sqrt{Q}$$

is the minimal incoming velocity of incident WIMPs that can deposit the recoil energy Q in the detector,

$$\mathcal{A} \equiv \frac{\rho_0 \sigma_0}{2m_\chi m_{r,N}^2} \quad \alpha \equiv \sqrt{\frac{m_N}{2m_{r,N}^2}} \quad m_{r,N} = \frac{m_\chi m_N}{m_\chi + m_N}$$

ρ_0 : WIMP density near the Earth

σ_0 : total cross section ignoring the form factor suppression

$F(Q)$: elastic nuclear form factor

$f_1(v)$: one-dimensional velocity distribution of halo WIMPs



Reconstruction of the 1-D WIMP velocity distribution



With measured recoil energies

Reconstruction of the 1-D WIMP velocity distribution

- Normalized one-dimensional WIMP velocity distribution function

$$f_1(v) = \mathcal{N} \left\{ -2Q \cdot \frac{d}{dQ} \left[\frac{1}{F^2(Q)} \left(\frac{dR}{dQ} \right) \right] \right\}_{Q=v^2/\alpha^2}$$

$$\mathcal{N} = \frac{2}{\alpha} \left\{ \int_0^\infty \frac{1}{\sqrt{Q}} \left[\frac{1}{F^2(Q)} \left(\frac{dR}{dQ} \right) \right] dQ \right\}^{-1}$$

- Moments of the velocity distribution function

$$\langle v^n \rangle = \mathcal{N}(Q_{\text{thre}}) \left(\frac{\alpha^{n+1}}{2} \right) \left[\frac{2Q_{\text{thre}}^{(n+1)/2}}{F^2(Q_{\text{thre}})} \left(\frac{dR}{dQ} \right)_{Q=Q_{\text{thre}}} + (n+1)I_n(Q_{\text{thre}}) \right]$$

$$\mathcal{N}(Q_{\text{thre}}) = \frac{2}{\alpha} \left[\frac{2Q_{\text{thre}}^{1/2}}{F^2(Q_{\text{thre}})} \left(\frac{dR}{dQ} \right)_{Q=Q_{\text{thre}}} + I_0(Q_{\text{thre}}) \right]^{-1}$$

$$I_n(Q_{\text{thre}}) = \int_{Q_{\text{thre}}}^\infty Q^{(n-1)/2} \left[\frac{1}{F^2(Q)} \left(\frac{dR}{dQ} \right) \right] dQ$$



Reconstruction of the 1-D WIMP velocity distribution

- **Ansatz:** the **measured** recoil spectrum in the n th Q -bin

$$\left(\frac{dR}{dQ}\right)_{\text{expt, } Q \simeq Q_n} \equiv r_n e^{k_n(Q - Q_{s,n})} \quad r_n \equiv \frac{N_n}{b_n}$$



Reconstruction of the 1-D WIMP velocity distribution

- **Ansatz:** the measured recoil spectrum in the n th Q -bin

$$\left(\frac{dR}{dQ}\right)_{\text{expt, } Q \simeq Q_n} \equiv r_n e^{k_n(Q - Q_{s,n})} \quad r_n \equiv \frac{N_n}{b_n}$$

- Logarithmic slope and shifted point in the n th Q -bin

$$\overline{Q - Q_n}|_n \equiv \frac{1}{N_n} \sum_{i=1}^{N_n} (Q_{n,i} - Q_n) = \left(\frac{b_n}{2}\right) \coth\left(\frac{k_n b_n}{2}\right) - \frac{1}{k_n}$$

$$Q_{s,n} = Q_n + \frac{1}{k_n} \ln \left[\frac{\sinh(k_n b_n / 2)}{k_n b_n / 2} \right]$$

Reconstruction of the 1-D WIMP velocity distribution

- **Ansatz:** the measured recoil spectrum in the n th Q -bin

$$\left(\frac{dR}{dQ}\right)_{\text{expt}, Q \simeq Q_n} \equiv r_n e^{k_n(Q - Q_{s,n})} \quad r_n \equiv \frac{N_n}{b_n}$$

- Logarithmic slope and shifted point in the n th Q -bin

$$\overline{Q - Q_n}|_n \equiv \frac{1}{N_n} \sum_{i=1}^{N_n} (Q_{n,i} - Q_n) = \left(\frac{b_n}{2}\right) \coth\left(\frac{k_n b_n}{2}\right) - \frac{1}{k_n}$$

$$Q_{s,n} = Q_n + \frac{1}{k_n} \ln \left[\frac{\sinh(k_n b_n / 2)}{k_n b_n / 2} \right]$$

- Reconstructing the one-dimensional WIMP velocity distribution

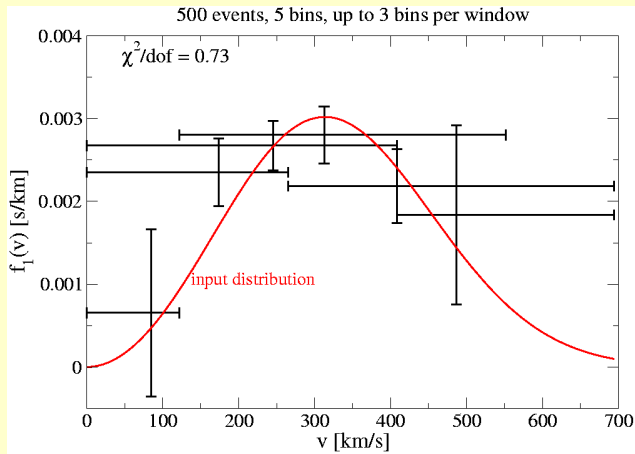
$$f_1(v_{s,n}) = \mathcal{N} \left[\frac{2Q_{s,n} r_n}{F^2(Q_{s,n})} \right] \left[\frac{d}{dQ} \ln F^2(Q) \Big|_{Q=Q_{s,n}} - k_n \right]$$

$$\mathcal{N} = \frac{2}{\alpha} \left[\sum_a \frac{1}{\sqrt{Q_a} F^2(Q_a)} \right]^{-1} \quad v_{s,n} = \alpha \sqrt{Q_{s,n}}$$

[M. Drees and CLS, JCAP 0706, 011 (2007)]

Reconstruction of the 1-D WIMP velocity distribution

- Reconstructed $f_{1,rec}(v_{s,n})$
(^{76}Ge , 500 events, 5 bins, up to 3 bins per window)



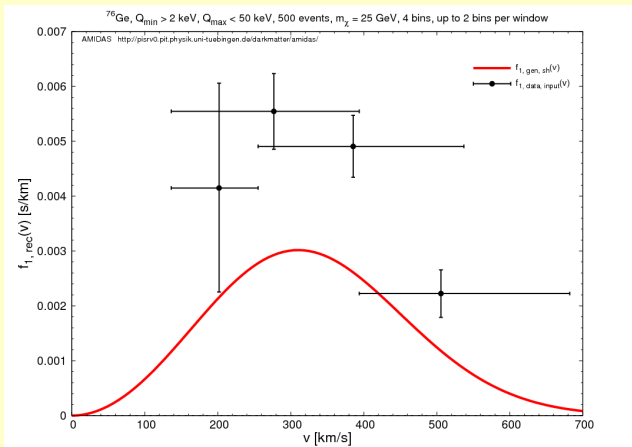
[M. Drees and CLS, JCAP 0706, 011 (2007)]



With a non-negligible threshold energy

Reconstruction of $f_1(v)$ with a non-negligible threshold energy

- Reconstructed $f_{1,\text{rec}}(v_{S,n})$ with a non-negligible threshold energy
(^{76}Ge , 2 - 50 keV, 500 events, $m_\chi = 25$ GeV)



[CLS, IJMPD 24, 1550090 (2015)]

Reconstruction of $f_1(v)$ with a non-negligible threshold energy

- Modification of the renormalization constant

$$\mathcal{N} = \frac{2}{\alpha} \left[\tilde{f}_{1,\text{rec}}(v_{\min}^*) Q_{\min}^{1/2} + \frac{2Q_{\min}^{1/2}}{F^2(Q_{\min})} \left(\frac{dR}{dQ} \right)_{\text{expt}, Q=Q_{\min}} + I_0(Q_{\min}, Q_{\max}^*) \right]^{-1}$$

where

$$\tilde{f}_{1,\text{rec}}(v_{\min}^*) \equiv \left[\frac{2Q_{\min} r(Q_{\min})}{F^2(Q_{\min})} \right] \left[\left. \frac{d}{dQ} \ln F^2(Q) \right|_{Q=Q_{\min}} - k_1 \right]$$

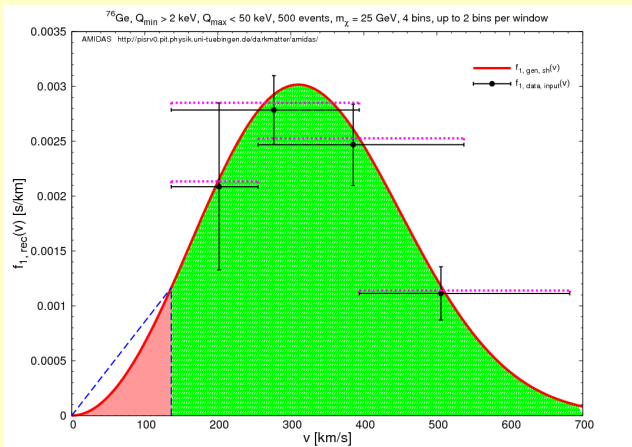
$$\left(\frac{dR}{dQ} \right)_{\text{expt}, Q=Q_{\min}} = r_1 e^{k_1(Q_{\min} - Q_{s,1})} \equiv r(Q_{\min})$$

$$I_n(Q_{\min}, Q_{\max}^*) = \int_{Q_{\min}}^{Q_{\max}^*} Q^{(n-1)/2} \left[\frac{1}{F^2(Q)} \left(\frac{dR}{dQ} \right) \right] dQ \rightarrow \sum_a \frac{Q_a^{(n-1)/2}}{F^2(Q_a)}$$

$$Q_{\max}^* \equiv \min \left(Q_{\max}, Q_{\max,\text{kin}} = \frac{v_{\max}^2}{\alpha^2} \right)$$

Reconstruction of $f_1(v)$ with a non-negligible threshold energy

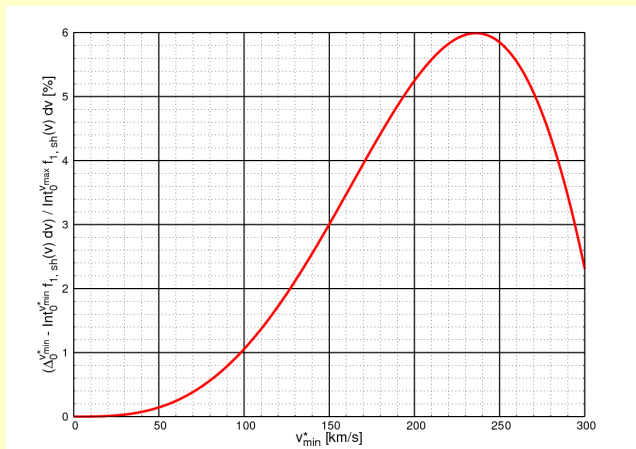
- Reconstructed $f_{1,rec}(v_{S,n})$ with the input WIMP mass
(^{76}Ge , 2 - 50 keV, 500 events, $m_\chi = 25$ GeV)



[CLS, IJMPD 24, 1550090 (2015)]

Reconstruction of $f_1(v)$ with a non-negligible threshold energy

- Theoretical bias estimate of $\left[\Delta_0^{v_{\min}^*} - \int_0^{v_{\min}^*} f_1(v) dv \right] / \int_0^{v_{\max}} f_1(v) dv$



[CLS, IJMPD 24, 1550090 (2015)]



With a model of the WIMP velocity distribution – Bayesian analysis



Bayesian reconstruction of $f_1(v)$

□ Bayesian analysis

$$p(\Theta|\text{data}) = \frac{p(\text{data}|\Theta)}{p(\text{data})} \cdot p(\Theta)$$

- Θ : $\{a_1, a_2, \dots, a_{N_{\text{Bayesian}}}\}$, a specified (combination of the) value(s) of the fitting parameter(s)
- $p(\Theta)$: **prior probability**, our degree of belief about Θ being the true value(s) of fitting parameter(s), often given in form of the **(multiplication of the) probability distribution(s)** of the fitting parameter(s)
- $p(\text{data}|\Theta)$: the probability of the observed result, once the specified (combination of the) value(s) of the fitting parameter(s) happens, usually be described by the **“likelihood” function of Θ , $\mathcal{L}(\Theta)$** .
- $p(\text{data})$: **evidence**, the total probability of obtaining the particular set of data
- $p(\Theta|\text{data})$: **posterior probability density function for Θ** , the probability of that the specified (combination of the) value(s) of the fitting parameter(s) happens, given the observed result



Bayesian reconstruction of $f_1(v)$

- Probability distribution functions for $p(\Theta)$
 - **Without** prior knowledge about the fitting parameter
 - ⇨ Flat-distributed

$$p_i(\mathbf{a}_i) = 1 \quad \text{for } a_{i,\min} \leq \mathbf{a}_i \leq a_{i,\max}$$

- **With** prior knowledge about the fitting parameter
 - ⇨ Around a theoretical predicted/estimated or experimental measured value $\mu_{a,i}$
 - ⇨ With (statistical) uncertainties $\sigma_{a,i}$
 - ⇨ Gaussian-distributed

$$p_i(\mathbf{a}_i; \mu_{a,i}, \sigma_{a,i}) = \frac{1}{\sqrt{2\pi} \sigma_{a,i}} e^{-(\mathbf{a}_i - \mu_{a,i})^2 / 2\sigma_{a,i}^2}$$

Bayesian reconstruction of $f_1(v)$

□ Likelihood function for $p(\text{data}|\Theta)$

- Theoretical one-dimensional WIMP velocity distribution function:
 $f_{1,\text{th}}(v; \mathbf{a}_1, \mathbf{a}_2, \dots, \mathbf{a}_{N_{\text{Bayesian}}})$
- Assuming that the reconstructed data points are **Gaussian**-distributed around the theoretical predictions

$$\mathcal{L}\left(f_{1,\text{rec}}(v_{s,\mu}), \mu = 1, 2, \dots, W; \mathbf{a}_i, i = 1, 2, \dots, N_{\text{Bayesian}}\right)$$

$$\equiv \prod_{\mu=1}^W \text{Gau}\left(v_{s,\mu}, f_{1,\text{rec}}(v_{s,\mu}), \sigma_{f_{1,s,\mu}}; \mathbf{a}_1, \mathbf{a}_2, \dots, \mathbf{a}_{N_{\text{Bayesian}}}\right)$$

with

$$\text{Gau}\left(v_{s,\mu}, f_{1,\text{rec}}(v_{s,\mu}), \sigma_{f_{1,s,\mu}}; \mathbf{a}_1, \mathbf{a}_2, \dots, \mathbf{a}_{N_{\text{Bayesian}}}\right)$$

$$\equiv \frac{1}{\sqrt{2\pi} \sigma_{f_{1,s,\mu}}} e^{-\left[f_{1,\text{rec}}(v_{s,\mu}) - f_{1,\text{th}}(v_{s,\mu}; \mathbf{a}_1, \mathbf{a}_2, \dots, \mathbf{a}_{N_{\text{Bayesian}}})\right]^2 / 2\sigma_{f_{1,s,\mu}}^2}$$



Bayesian reconstruction of $f_1(v)$

□ Input and fitting one-dimensional WIMP velocity distribution functions

➤ “One-parameter” shifted Maxwellian velocity distribution

$$f_{1,\text{sh},v_0}(v) = \frac{1}{\sqrt{\pi}} \left(\frac{v}{v_0 v_e} \right) \left[e^{-(v-v_e)^2/v_0^2} - e^{-(v+v_e)^2/v_0^2} \right] \quad v_e = 1.05 v_0$$

➤ Shifted Maxwellian velocity distribution

$$f_{1,\text{sh}}(v) = \frac{1}{\sqrt{\pi}} \left(\frac{v}{v_0 v_e} \right) \left[e^{-(v-v_e)^2/v_0^2} - e^{-(v+v_e)^2/v_0^2} \right]$$

➤ “Variatied” shifted Maxwellian velocity distribution

$$f_{1,\text{sh},\Delta v}(v) = \frac{1}{\sqrt{\pi}} \left[\frac{v}{v_0 (v_0 + \Delta v)} \right] \left\{ e^{-[v-(v_0+\Delta v)]^2/v_0^2} - e^{-[v+(v_0+\Delta v)]^2/v_0^2} \right\}$$



Bayesian reconstruction of $f_1(v)$

□ Input and fitting one-dimensional WIMP velocity distribution functions

➤ “One-parameter” shifted Maxwellian velocity distribution

$$f_{1,\text{sh},v_0}(v) = \frac{1}{\sqrt{\pi}} \left(\frac{v}{v_0 v_e} \right) \left[e^{-(v-v_e)^2/v_0^2} - e^{-(v+v_e)^2/v_0^2} \right] \quad v_e = 1.05 v_0$$

➤ Shifted Maxwellian velocity distribution

$$f_{1,\text{sh}}(v) = \frac{1}{\sqrt{\pi}} \left(\frac{v}{v_0 v_e} \right) \left[e^{-(v-v_e)^2/v_0^2} - e^{-(v+v_e)^2/v_0^2} \right]$$

➤ “Variatied” shifted Maxwellian velocity distribution

$$f_{1,\text{sh},\Delta v}(v) = \frac{1}{\sqrt{\pi}} \left[\frac{v}{v_0 (v_0 + \Delta v)} \right] \left\{ e^{-[v-(v_0+\Delta v)]^2/v_0^2} - e^{-[v+(v_0+\Delta v)]^2/v_0^2} \right\}$$

➤ Simple Maxwellian velocity distribution

$$f_{1,\text{Gau}}(v) = \frac{4}{\sqrt{\pi}} \left(\frac{v^2}{v_0^3} \right) e^{-v^2/v_0^2}$$

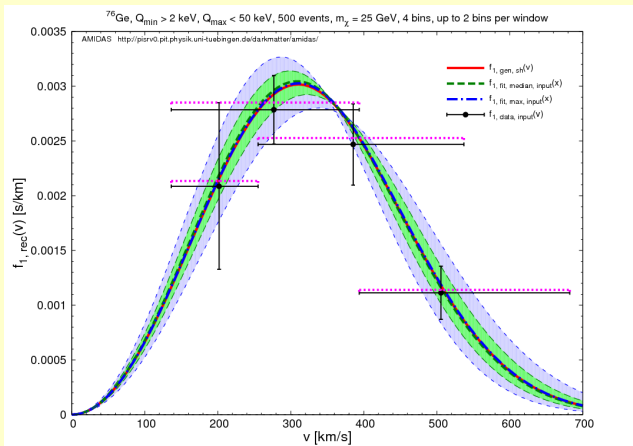
➤ “Modified” simple Maxwellian velocity distribution

$$f_{1,\text{Gau},k}(v) = \frac{v^2}{N_{f,k}} \left(e^{-v^2/kv_0^2} - e^{-v_{\text{max}}^2/kv_0^2} \right)^k \quad \text{for } v \leq v_{\text{max}}$$

Bayesian reconstruction of $f_1(v)$

- Reconstructed $f_{1,\text{Bayesian}}(v)$ with the input WIMP mass

(^{76}Ge , $Q_{\text{min}} > 2$ keV, $Q_{\text{max}} < 50$ keV, 500 events, $m_\chi = 25$ GeV, $f_{1,\text{sh},v_0}(v) \Rightarrow f_{1,\text{sh},v_0}(v)$, flat-dist.)



[CLS, IJMPD 24, 1550090 (2015)]



Determination of the WIMP mass



Determination of the WIMP mass

- Estimating the moments of the WIMP velocity distribution

$$\langle v^n \rangle = \alpha^n \left[\frac{2Q_{\min}^{1/2} r_{\min}}{F^2(Q_{\min})} + I_0 \right]^{-1} \left[\frac{2Q_{\min}^{(n+1)/2} r_{\min}}{F^2(Q_{\min})} + (n+1)I_n \right]$$

$$I_n = \sum_a \frac{Q_a^{(n-1)/2}}{F^2(Q_a)}$$

$$r_{\min} = \left(\frac{dR}{dQ} \right)_{\text{expt}, Q=Q_{\min}} = r_1 e^{k_1(Q_{\min} - Q_{s,1})}$$

[M. Drees and CLS, JCAP 0706, 011 (2007)]



Determination of the WIMP mass

- Estimating the moments of the WIMP velocity distribution

$$\langle v^n \rangle = \alpha^n \left[\frac{2Q_{\min}^{1/2} r_{\min}}{F^2(Q_{\min})} + I_0 \right]^{-1} \left[\frac{2Q_{\min}^{(n+1)/2} r_{\min}}{F^2(Q_{\min})} + (n+1)I_n \right]$$

$$I_n = \sum_a \frac{Q_a^{(n-1)/2}}{F^2(Q_a)} \quad r_{\min} = \left(\frac{dR}{dQ} \right)_{\text{expt}, Q=Q_{\min}} = r_1 e^{k_1(Q_{\min} - Q_{s,1})}$$

[M. Drees and CLS, JCAP 0706, 011 (2007)]

- Determining the WIMP mass

$$m_X |_{\langle v^n \rangle} = \frac{\sqrt{m_X m_Y} - m_X \mathcal{R}_n}{\mathcal{R}_n - \sqrt{m_X / m_Y}}$$

$$\mathcal{R}_n = \left[\frac{2Q_{\min,X}^{(n+1)/2} r_{\min,X} / F_X^2(Q_{\min,X}) + (n+1)I_{n,X}}{2Q_{\min,X}^{1/2} r_{\min,X} / F_X^2(Q_{\min,X}) + I_{0,X}} \right]^{1/n} (X \rightarrow Y)^{-1} \quad (n \neq 0)$$

[CLS and M. Drees, arXiv:0710.4296]

Determination of the WIMP mass

- Estimating the moments of the WIMP velocity distribution

$$\langle v^n \rangle = \alpha^n \left[\frac{2Q_{\min}^{1/2} r_{\min}}{F^2(Q_{\min})} + I_0 \right]^{-1} \left[\frac{2Q_{\min}^{(n+1)/2} r_{\min}}{F^2(Q_{\min})} + (n+1)I_n \right]$$

$$I_n = \sum_a \frac{Q_a^{(n-1)/2}}{F^2(Q_a)} \quad r_{\min} = \left(\frac{dR}{dQ} \right)_{\text{expt}, Q=Q_{\min}} = r_1 e^{k_1(Q_{\min} - Q_{s,1})}$$

[M. Drees and CLS, JCAP 0706, 011 (2007)]

- Determining the WIMP mass

$$m_X |_{\langle v^n \rangle} = \frac{\sqrt{m_X m_Y} - m_X \mathcal{R}_n}{\mathcal{R}_n - \sqrt{m_X / m_Y}}$$

$$\mathcal{R}_n = \left[\frac{2Q_{\min,X}^{(n+1)/2} r_{\min,X} / F_X^2(Q_{\min,X}) + (n+1)I_{n,X}}{2Q_{\min,X}^{1/2} r_{\min,X} / F_X^2(Q_{\min,X}) + I_{0,X}} \right]^{1/n} (X \rightarrow Y)^{-1} \quad (n \neq 0)$$

[CLS and M. Drees, arXiv:0710.4296]

- Assuming a dominant SI scalar WIMP-nucleus interaction

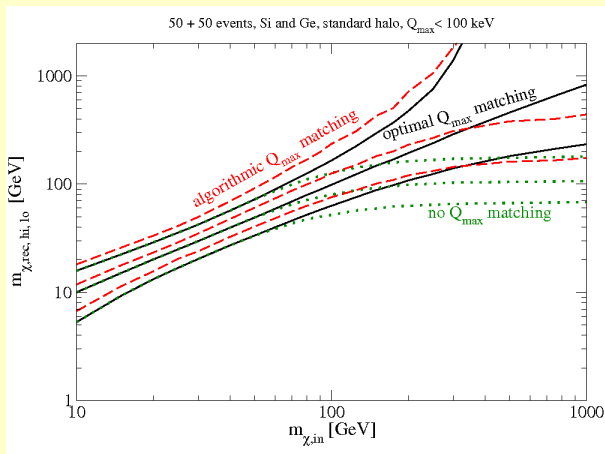
$$m_X |_{\sigma} = \frac{(m_X / m_Y)^{5/2} m_Y - m_X \mathcal{R}_{\sigma}}{\mathcal{R}_{\sigma} - (m_X / m_Y)^{5/2}}$$

$$\mathcal{R}_{\sigma} = \frac{\mathcal{E}_Y}{\mathcal{E}_X} \left[\frac{2Q_{\min,X}^{1/2} r_{\min,X} / F_X^2(Q_{\min,X}) + I_{0,X}}{2Q_{\min,Y}^{1/2} r_{\min,Y} / F_Y^2(Q_{\min,Y}) + I_{0,Y}} \right]$$

[M. Drees and CLS, JCAP 0806, 012 (2008)]

Determination of the WIMP mass

- Reconstructed $m_{\chi, \text{rec}}$
($^{28}\text{Si} + ^{76}\text{Ge}$, $Q_{\text{max}} < 100 \text{ keV}$, $2 \times 50 \text{ events}$)



[M. Drees and CLS, JCAP 0806, 012 (2008)]



Determinations of the WIMP-nucleon couplings

- └ Determinations of the WIMP-nucleon couplings
 - └ Estimation of the SI scalar WIMP-nucleon coupling



Estimation of the SI scalar WIMP-nucleon coupling



Estimation of the SI scalar WIMP-nucleon coupling

- Spin-independent (SI) scalar WIMP-nucleus cross section

$$\sigma_0^{\text{SI}} = \left(\frac{4}{\pi}\right) m_{r,N}^2 [Z f_p + (A - Z) f_n]^2 \simeq \left(\frac{4}{\pi}\right) m_{r,N}^2 A^2 |f_p|^2 = A^2 \left(\frac{m_{r,N}}{m_{r,p}}\right)^2 \sigma_{\chi p}^{\text{SI}}$$

$$\sigma_{\chi p}^{\text{SI}} = \left(\frac{4}{\pi}\right) m_{r,p}^2 |f_p|^2$$

$f_{(p,n)}$: effective SI scalar WIMP-proton/neutron couplings

Estimation of the SI scalar WIMP-nucleon coupling

- Spin-independent (SI) scalar WIMP-nucleus cross section

$$\sigma_0^{\text{SI}} = \left(\frac{4}{\pi}\right) m_{r,N}^2 [Z f_p + (A - Z) f_n]^2 \simeq \left(\frac{4}{\pi}\right) m_{r,N}^2 A^2 |f_p|^2 = A^2 \left(\frac{m_{r,N}}{m_{r,p}}\right)^2 \sigma_{\chi p}^{\text{SI}}$$

$$\sigma_{\chi p}^{\text{SI}} = \left(\frac{4}{\pi}\right) m_{r,p}^2 |f_p|^2$$

$f_{(p,n)}$: effective SI scalar WIMP-proton/neutron couplings

- Rewriting the integral over $f_1(v)/v$

$$\left(\frac{dR}{dQ}\right)_{\text{expt, } Q=Q_{\min}} = \frac{\mathcal{E} \rho_0 A^2}{2m_\chi m_{r,p}^2} \left[\left(\frac{4}{\pi}\right) m_{r,p}^2 |f_p|^2 \right] F^2(Q_{\min}) \left\{ m_{r,N} \sqrt{\frac{2}{m_N}} \left[\frac{2Q_{\min}^{1/2} r_{\min}}{F^2(Q_{\min})} + l_0 \right]^{-1} \left[\frac{2r_{\min}}{F^2(Q_{\min})} \right] \right\}$$

- └ Determinations of the WIMP-nucleon couplings
 - └ Estimation of the SI scalar WIMP-nucleon coupling



Estimation of the SI scalar WIMP-nucleon coupling

- Spin-independent (SI) scalar WIMP-nucleus cross section

$$\sigma_0^{\text{SI}} = \left(\frac{4}{\pi}\right) m_{r,N}^2 [Z f_p + (A - Z) f_n]^2 \simeq \left(\frac{4}{\pi}\right) m_{r,N}^2 A^2 |f_p|^2 = A^2 \left(\frac{m_{r,N}}{m_{r,p}}\right)^2 \sigma_{\chi p}^{\text{SI}}$$

$$\sigma_{\chi p}^{\text{SI}} = \left(\frac{4}{\pi}\right) m_{r,p}^2 |f_p|^2$$

$f_{(p,n)}$: effective SI scalar WIMP-proton/neutron couplings

- Rewriting the integral over $f_1(v)/v$

$$\left(\frac{dR}{dQ}\right)_{\text{expt, } Q=Q_{\min}} = \frac{\mathcal{E} \rho_0 A^2}{2m_\chi m_{r,p}^2} \left[\left(\frac{4}{\pi}\right) m_{r,p}^2 |f_p|^2 \right] F^2(Q_{\min}) \left\{ m_{r,N} \sqrt{\frac{2}{m_N}} \left[\frac{2Q_{\min}^{1/2} r_{\min}}{F^2(Q_{\min})} + l_0 \right]^{-1} \left[\frac{2r_{\min}}{F^2(Q_{\min})} \right] \right\}$$

- Estimating the SI scalar WIMP-nucleon coupling

$$|f_p|^2 = \frac{1}{\rho_0} \left[\frac{\pi}{4\sqrt{2}} \left(\frac{1}{\mathcal{E}_Z A_Z^2 \sqrt{m_Z}} \right) \right] \left[\frac{2Q_{\min,Z}^{1/2} r_{\min,Z}}{F_Z^2(Q_{\min,Z})} + l_{0,Z} \right] (m_\chi + m_Z)$$

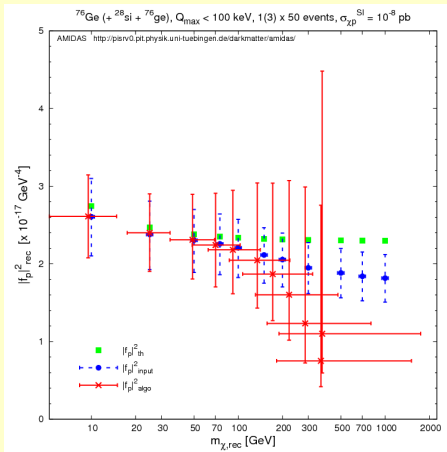
[M. Drees and CLS, PoS IDM2008, 110 (2008)]

- └ Determinations of the WIMP-nucleon couplings
 - └ Estimation of the SI scalar WIMP-nucleon coupling



Estimation of the SI scalar WIMP-nucleon coupling

- Reconstructed $|f_p|_{\text{rec}}^2$ vs. reconstructed $m_{\chi, \text{rec}}$
 (^{76}Ge (+ ^{28}Si + ^{76}Ge), $Q_{\text{max}} < 100$ keV, $1(3) \times 50$ events)



[CLS, arXiv:1103.0481]



Determination of the ratio of SD WIMP-nucleon couplings

- Spin-dependent (SD) axial-vector WIMP-nucleus cross section

$$\sigma_0^{\text{SD}} = \left(\frac{32}{\pi}\right) G_F^2 m_{r,N}^2 \left(\frac{J+1}{J}\right) [\langle S_p \rangle a_p + \langle S_n \rangle a_n]^2$$

$$\sigma_{\chi_{p/n}}^{\text{SD}} = \left(\frac{32}{\pi}\right) G_F^2 m_{r,p/n}^2 \cdot \left(\frac{3}{4}\right) a_{p/n}^2$$

J : total nuclear spin

$\langle S_{(p,n)} \rangle$: expectation values of the proton/neutron group spin

$a_{(p,n)}$: effective SD axial-vector WIMP-proton/neutron couplings

Determination of the ratio of SD WIMP-nucleon couplings

- Spin-dependent (SD) axial-vector WIMP-nucleus cross section

$$\sigma_0^{\text{SD}} = \left(\frac{32}{\pi}\right) G_F^2 m_{r,N}^2 \left(\frac{J+1}{J}\right) [\langle S_p \rangle a_p + \langle S_n \rangle a_n]^2$$

$$\sigma_{\chi_{p/n}}^{\text{SD}} = \left(\frac{32}{\pi}\right) G_F^2 m_{r,p/n}^2 \cdot \left(\frac{3}{4}\right) a_{p/n}^2$$

J : total nuclear spin

$\langle S_{(p,n)} \rangle$: expectation values of the proton/neutron group spin

$a_{(p,n)}$: effective SD axial-vector WIMP-proton/neutron couplings

- Determining the ratio of two SD axial-vector WIMP-nucleon couplings

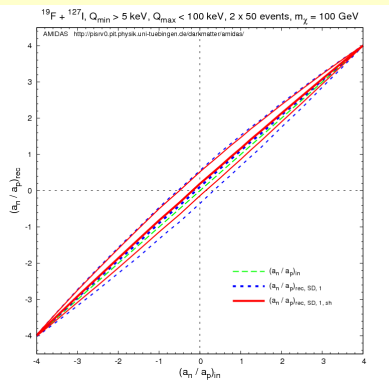
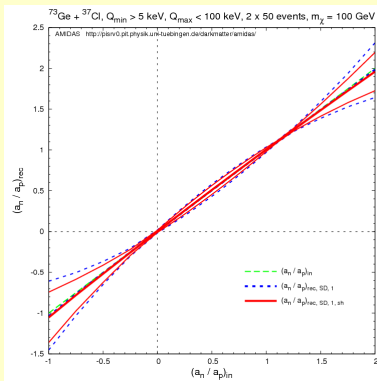
$$\left(\frac{a_n}{a_p}\right)_{\pm,n}^{\text{SD}} = -\frac{\langle S_p \rangle_X \pm \langle S_p \rangle_Y \mathcal{R}_{J,n}}{\langle S_n \rangle_X \pm \langle S_n \rangle_Y \mathcal{R}_{J,n}}$$

$$\mathcal{R}_{J,n} \equiv \left[\left(\frac{J_X}{J_X+1}\right) \left(\frac{J_Y+1}{J_Y}\right) \frac{\mathcal{R}_\sigma}{\mathcal{R}_n} \right]^{1/2} \quad (n \neq 0)$$

Determination of the ratio of SD WIMP-nucleon couplings

□ Reconstructed $(a_n/a_p)^{SD}_{rec,1}$

($^{73}\text{Ge} + ^{37}\text{Cl}$ and $^{19}\text{F} + ^{127}\text{I}$, $Q_{min} > 5 \text{ keV}$, $Q_{max} < 100 \text{ keV}$, $2 \times 50 \text{ events}$,
 $m_\chi = 100 \text{ GeV}$)





Determinations of ratios of WIMP-nucleon cross sections

- Differential rate for **combined SI and SD** cross sections

$$\left(\frac{dR}{dQ}\right)_{\text{expt, } Q=Q_{\min}} = \mathcal{E} \left(\frac{\rho_0 \sigma_0^{\text{SI}}}{2m_\chi m_{r,N}^2} \right) \left[F_{\text{SI}}^2(Q) + \left(\frac{\sigma_{\chi p}^{\text{SD}}}{\sigma_{\chi p}^{\text{SI}}} \right) c_p F_{\text{SD}}^2(Q) \right] \int_{v_{\min}}^{v_{\max}} \left[\frac{f_1(v)}{v} \right] dv$$

$$c_p \equiv \frac{4}{3} \left(\frac{J+1}{J} \right) \left[\frac{\langle S_p \rangle + (a_n/a_p) \langle S_n \rangle}{A} \right]^2$$



Determinations of ratios of WIMP-nucleon cross sections

- Differential rate for **combined SI and SD** cross sections

$$\left(\frac{dR}{dQ}\right)_{\text{expt}, Q=Q_{\min}} = \mathcal{E} \left(\frac{\rho_0 \sigma_0^{\text{SI}}}{2m_\chi m_{r,N}^2} \right) \left[F_{\text{SI}}^2(Q) + \left(\frac{\sigma_{\chi\text{p}}^{\text{SD}}}{\sigma_{\chi\text{p}}^{\text{SI}}} \right) c_p F_{\text{SD}}^2(Q) \right] \int_{v_{\min}}^{v_{\max}} \left[\frac{f_1(v)}{v} \right] dv$$

$$c_p \equiv \frac{4}{3} \left(\frac{J+1}{J} \right) \left[\frac{\langle S_p \rangle + (a_n/a_p) \langle S_n \rangle}{A} \right]^2$$

- Determining the **ratio of two WIMP-proton cross sections**

$$\frac{\sigma_{\chi\text{p}}^{\text{SD}}}{\sigma_{\chi\text{p}}^{\text{SI}}} = \frac{F_{\text{SI},Y}^2(Q_{\min,Y}) \mathcal{R}_{m,XY} - F_{\text{SI},X}^2(Q_{\min,X})}{c_{p,X} F_{\text{SD},X}^2(Q_{\min,X}) - c_{p,Y} F_{\text{SD},Y}^2(Q_{\min,Y}) \mathcal{R}_{m,XY}}$$

$$\mathcal{R}_{m,XY} \equiv \left(\frac{r_{\min,X}}{\mathcal{E}_X} \right) \left(\frac{\mathcal{E}_Y}{r_{\min,Y}} \right) \left(\frac{m_Y}{m_X} \right)^2$$

Determinations of ratios of WIMP-nucleon cross sections

- Differential rate for **combined SI and SD** cross sections

$$\left(\frac{dR}{dQ}\right)_{\text{expt, } Q=Q_{\min}} = \mathcal{E} \left(\frac{\rho_0 \sigma_0^{\text{SI}}}{2m_\chi m_{r,N}^2} \right) \left[F_{\text{SI}}^2(Q) + \left(\frac{\sigma_{\chi\text{P}}^{\text{SD}}}{\sigma_{\chi\text{P}}^{\text{SI}}} \right) c_p F_{\text{SD}}^2(Q) \right] \int_{v_{\min}}^{v_{\max}} \left[\frac{f_1(v)}{v} \right] dv$$

$$c_p \equiv \frac{4}{3} \left(\frac{J+1}{J} \right) \left[\frac{\langle S_p \rangle + (a_n/a_p) \langle S_n \rangle}{A} \right]^2$$

- Determining the **ratio of two WIMP-proton cross sections**

$$\frac{\sigma_{\chi\text{P}}^{\text{SD}}}{\sigma_{\chi\text{P}}^{\text{SI}}} = \frac{F_{\text{SI},Y}^2(Q_{\min},Y) \mathcal{R}_{m,XY} - F_{\text{SI},X}^2(Q_{\min},X)}{c_{p,X} F_{\text{SD},X}^2(Q_{\min},X) - c_{p,Y} F_{\text{SD},Y}^2(Q_{\min},Y) \mathcal{R}_{m,XY}}$$

$$\mathcal{R}_{m,XY} \equiv \left(\frac{r_{\min,X}}{\mathcal{E}_X} \right) \left(\frac{\mathcal{E}_Y}{r_{\min,Y}} \right) \left(\frac{m_Y}{m_X} \right)^2$$

- Determining the **ratio of two SD axial-vector WIMP-nucleon couplings**

$$\left(\frac{a_n}{a_p} \right)_{\pm}^{\text{SI+SD}} = \frac{- \left(c_{p,X} s_{n/p,X} - c_{p,Y} s_{n/p,Y} \right) \pm \sqrt{c_{p,X} c_{p,Y}} \left| s_{n/p,X} - s_{n/p,Y} \right|}{c_{p,X} s_{n/p,X}^2 - c_{p,Y} s_{n/p,Y}^2}$$

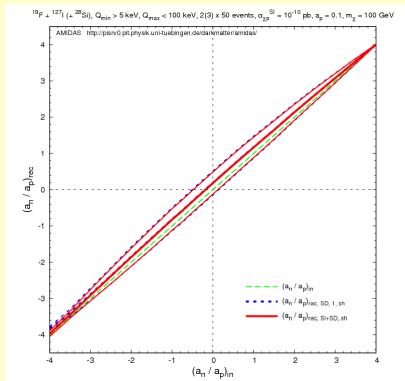
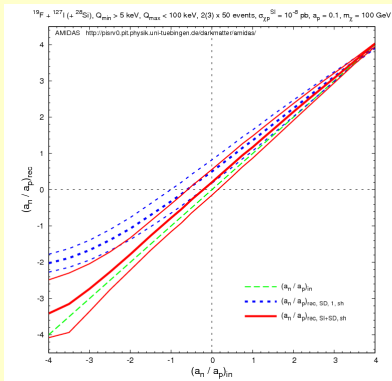
$$c_{p,X} \equiv \frac{4}{3} \left(\frac{J_X + 1}{J_X} \right) \left[\frac{\langle S_p \rangle_X}{A_X} \right]^2 \left[F_{\text{SI},Z}^2(Q_{\min},Z) \mathcal{R}_{m,YZ} - F_{\text{SI},Y}^2(Q_{\min},Y) \right] F_{\text{SD},X}^2(Q_{\min},X)$$

[M. Drees and CLS, arXiv:0903.3300]

Determinations of ratios of WIMP-nucleon cross sections

□ Reconstructed $(a_n/a_p)_{rec}^{SI+SD}$ vs. $(a_n/a_p)_{rec,1}^{SD}$

$(^{19}\text{F} + ^{127}\text{I} + ^{28}\text{Si}, Q_{min} > 5 \text{ keV}, Q_{max} < 100 \text{ keV}, 3 \times 50 \text{ events},$
 $\sigma_{\chi p}^{SI} = 10^{-8}/10^{-10} \text{ pb}, a_p = 0.1, m_\chi = 100 \text{ GeV})$

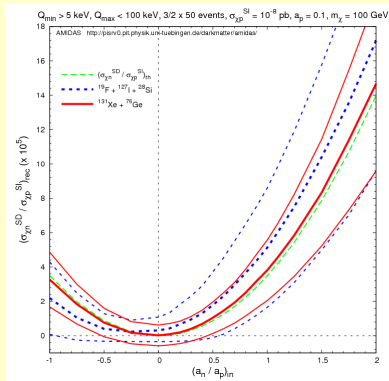
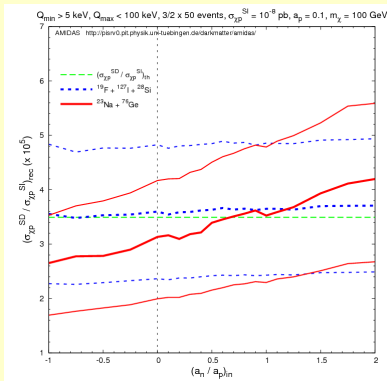


[CLS, JCAP 1107, 005 (2011)]

Determinations of ratios of WIMP-nucleon cross sections

□ Reconstructed $(\sigma_{\chi p}^{SD}/\sigma_{\chi p}^{SI})_{rec}$ and $(\sigma_{\chi n}^{SD}/\sigma_{\chi p}^{SI})_{rec}$

($^{19}\text{F} + ^{127}\text{I} + ^{28}\text{Si}$ vs. $^{23}\text{Na}/^{131}\text{Xe} + ^{76}\text{Ge}$, $Q_{min} > 5$ keV, $Q_{max} < 100$ keV,
 $\sigma_{\chi p}^{SI} = 10^{-8}$ pb, $a_p = 0.1$, $m_\chi = 100$ GeV, $3/2 \times 50$ events)



[CLS, JCAP 1107, 005 (2011)]



Summary



Summary

- Once two or more experiments with different target nuclei observe positive WIMP signals, we could reconstruct
 - WIMP mass m_χ
 - 1-D velocity distribution $f_1(v)$
 - SI WIMP-proton coupling $|f_p|^2$ (with an assumed ρ_0)
 - ratio between the SD WIMP-nucleon couplings a_n/a_p
 - ratios between the SD and SI WIMP-nucleon cross sections $\sigma_{\chi(p,n)}^{SD}/\sigma_{\chi p}^{SI}$



Summary

- Once two or more experiments with different target nuclei observe positive WIMP signals, we could reconstruct
 - WIMP mass m_χ
 - 1-D velocity distribution $f_1(v)$
 - SI WIMP-proton coupling $|f_p|^2$ (with an assumed ρ_0)
 - ratio between the SD WIMP-nucleon couplings a_n/a_p
 - ratios between the SD and SI WIMP-nucleon cross sections $\sigma_{\chi(p,n)}^{SD}/\sigma_{\chi p}^{SI}$
- With an assumed $f_{1,th}(v)$, one can fit $f_1(v)$ by using Bayesian analysis.



Summary

- Once two or more experiments with different target nuclei observe positive WIMP signals, we could reconstruct
 - WIMP mass m_χ
 - 1-D velocity distribution $f_1(v)$
 - SI WIMP-proton coupling $|f_p|^2$ (with an assumed ρ_0)
 - ratio between the SD WIMP-nucleon couplings a_n/a_p
 - ratios between the SD and SI WIMP-nucleon cross sections $\sigma_{\chi(p,n)}^{SD}/\sigma_{\chi p}^{SI}$
- With an assumed $f_{1,th}(v)$, one can fit $f_1(v)$ by using Bayesian analysis.
- For these analyses the local density, the velocity distribution, and the mass/couplings on nucleons of halo WIMPs are not required priorly.



Summary

- Once two or more experiments with different target nuclei observe positive WIMP signals, we could reconstruct
 - WIMP mass m_χ
 - 1-D velocity distribution $f_1(v)$
 - SI WIMP-proton coupling $|f_p|^2$ (with an assumed ρ_0)
 - ratio between the SD WIMP-nucleon couplings a_n/a_p
 - ratios between the SD and SI WIMP-nucleon cross sections $\sigma_{\chi(p,n)}^{SD}/\sigma_{\chi p}^{SI}$
- With an assumed $f_{1,th}(v)$, one can fit $f_1(v)$ by using Bayesian analysis.
- For these analyses the local density, the velocity distribution, and the mass/couplings on nucleons of halo WIMPs are not required priorly.
- For a WIMP mass of $\mathcal{O}(100 \text{ GeV})$, with only $\mathcal{O}(50)$ events from one experiment and less than $\sim 20\%$ unrejected backgrounds, these quantities could be estimated with statistical uncertainties of $10\% - 40\%$.



Thank you very much for your attention!